

# Basics of photometry

K. Nilsson

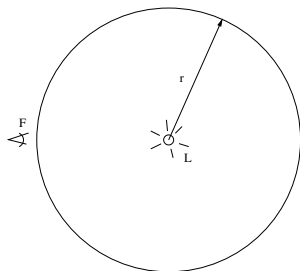
Finnish Centre for Astronomy with ESO (FINCA)

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# Luminosity and flux density

- An object has a total energy output = *Luminosity*  $L$  [erg s<sup>-1</sup>].
- The observer at distance  $r$  sees only part of this energy as flux density (flux)  $F$  [erg s<sup>-1</sup> cm<sup>-2</sup>].
- For an *isotropic source*

$$F = \frac{L}{4\pi r^2}.$$



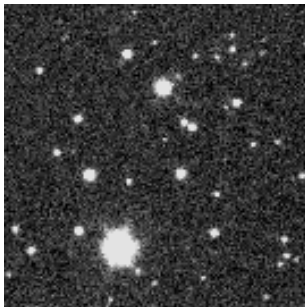
# Spectral flux

- Flux is the observed quantity, luminosity calculated (with a set of assumptions).
- Typically in astronomy one talks about flux when meaning *spectral* flux.
- Spectral flux has a unit of  $\text{erg s}^{-1} \text{cm}^{-2} \text{\AA}^{-1}$  ( $F_\lambda$ ) or  $\text{erg s}^{-1} \text{cm}^{-2} \text{Hz}^{-1}$  ( $F_\nu$ )
- Note that  $F_\lambda(\lambda_0) \neq F_\nu(\nu_0)$  although  $\lambda_0 = c/\nu_0$ , because  $d\lambda = -c/\nu^2 d\nu$ .

# The problem

CCD-cameras do not measure flux density  $F$  directly. They record something proportional to *energy* stored to the detector over the exposure time.

Light  $\rightarrow$  electrons  $\rightarrow$  voltage  $\rightarrow$  ADUs (“counts”)



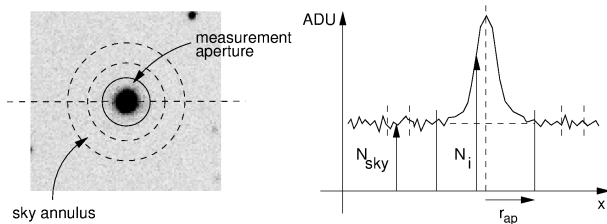
How to derive  $F$  from the pixel values on the CCD?

# Three problems

- 1) How to measure the counts from a target? Available techniques:
  - aperture photometry
  - PSF photometry
  - model fitting (e.g. galaxies).
- 2) How to correct the counts for
  - the atmosphere
  - instrumental response and artifacts.
- 3) Calibration: How to convert the measured counts to  $F$ ?
  - “Differential” vs. “absolute” calibration: In the former the calibration targets are on the same CCD frame as the target, in the latter they are on different CCD frames.

## 1. Counting the photons

# 1.1 Aperture photometry



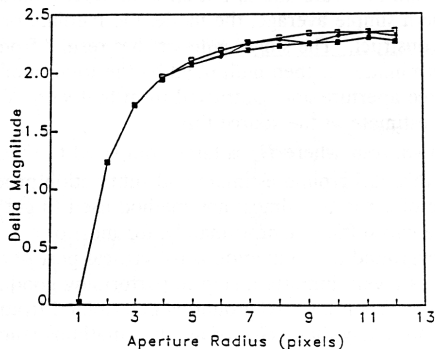
$$N = \sum_{i=1}^{n_{ap}} (N_i - N_{sky}) , \quad (1)$$

where  $N_i$  = counts in the  $i$ th pixel inside the aperture  
 $n_{ap}$  = number of pixels in the aperture ( $\sim \pi r_{ap}^2$ )  
 $N_{sky}$  = average sky brightness inside the aperture.

# How to choose aperture radius $r_{ap}$ ?

Guided by two facts:

- 1) *Growth curve*:  $N$  increases with increasing  $r_{ap}$ .



This means that one should use a very large aperture to include all target light, *but*

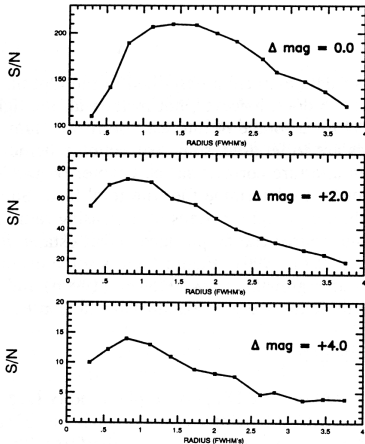


2) Maximum signal to noise is often achieved with a relatively small aperture, especially when the target is faint..

Signal to noise in aperture photometry:

$$S/N = \frac{N}{\sqrt{N + n_{ap}\sigma_{sky}^2(1 + \frac{n_{ap}}{n_{sky}})}} \quad (2)$$

$N$  = target counts [ $e^-$ ]  
 $\sigma_{sky}$  = rms sky noise [ $e^-$ ]  
 $n_{ap}$  = number of pixels in the aperture  
 $n_{sky}$  = number of pixels in the sky annulus.



# IRAF phot output (bright target)

```
c_e_20110206_16_1_1_1.f451.800 301.800 1 nullfile 0 \
451.193 304.201 -0.607 2.401 0.006 0.006 0 NoError \
69.56563 6.318079 2.231493 4689 36 0 NoError \
1. INDEF INDEF INDEF \
5.00 442399.4 78.80149 436917.6 10.899 0.001 0 NoError *\
7.50 496859.1 176.9911 484546.6 10.787 0.001 0 NoError *\
10.00 525029.1 314.4589 503153.6 10.746 0.001 0 NoError *\
12.50 547029.8 491.2861 512853.2 10.725 0.001 0 NoError *\
15.00 568272.6 707.141 519079.9 10.712 0.001 0 NoError *\
17.50 590793.5 962.5003 523836.5 10.702 0.001 0 NoError *\
20.00 615294.3 1256.851 527860.7 10.694 0.001 0 NoError *\
22.50 641928.2 1590.738 531267.5 10.687 0.001 0 NoError *\
25.00 671305. 1963.718 534697.8 10.680 0.001 0 NoError *\
27.50 703837.6 2376.03 538547.6 10.672 0.001 0 NoError *\
30.00 738392.7 2827.51 541695.2 10.666 0.001 0 NoError *\
35.00 813511.4 3848.458 545791. 10.657 0.001 0 NoError *\
40.00 899088. 5026.669 549404.6 10.650 0.002 0 NoError *\
45.00 994715.1 6362.06 552134.4 10.645 0.002 0 NoError *\
60.00 1345367. 11309.81 558592.4 10.632 0.003 0 NoError *\
70.00 1634051. 15393.9 563164.4 10.623 0.003 0 NoError *
rap N mag mag err
```

For bright targets aperture correction is not needed, the S/N is always good!

# IRAF phot output (faint target)

```
c_e_20110206_16_1_1_1.f196.200  484.200  2  nullfile  0  \  
195.618  483.560  -0.582  -0.640  0.070  0.068  0  NoError  \  
68.95222  6.23297  2.632393  4703  20  0  NoError  \  
1.  INDEF  INDEF  INDEF  \  
5.00  9369.124  78.84562  3932.542  16.013  0.019  0  NoError  *\  
7.50  16638.16  176.9813  4434.906  15.883  0.023  0  NoError  *\  
10.00  26254.11  314.3169  4581.261  15.848  0.029  0  NoError  *\  
12.50  38558.43  491.3904  4675.975  15.825  0.035  0  NoError  *\  
15.00  53700.28  706.8641  4960.433  15.761  0.040  0  NoError  *\  
17.50  71445.58  962.7172  5064.095  15.739  0.047  0  NoError  *\  
20.00  91834.17  1256.521  5194.255  15.711  0.053  0  NoError  *\  
22.50  114996.9  1591.084  5288.081  15.692  0.060  0  NoError  *\  
25.00  140757.4  1963.296  5383.831  15.672  0.067  0  NoError  *\  
27.50  169155.7  2376.665  5279.284  15.694  0.077  0  NoError  *\  
30.00  200351.3  2827.048  5420.033  15.665  0.085  0  NoError  *\  
35.00  270491.6  3848.048  5160.146  15.718  0.110  0  NoError  *\  
40.00  351731.4  5026.392  5150.483  15.720  0.134  0  NoError  *\  
45.00  444207.3  6362.046  5530.06  15.643  0.150  0  NoError  *\  
60.00  785383.  11310.65  5488.602  15.651  0.242  0  NoError  *\  
70.00  1067326.  15394.69  5828.265  15.586  0.298  0  NoError  *  
  
rap  N  mag  mag err
```

For faint targets the S/N quickly deteriorates at large radii.

# Choosing the aperture

Facts 1) and 2) give contradicting guidance:

- on one hand a large aperture should be used to obtain total  $N$
- on the other hand, best S/N is achieved with a small aperture, which does not yield total  $N$ .

Solution : Use *aperture correction*:

- Measure the target counts  $N$  inside the “best S/N” radius  $r_{ap}^{(1)}$  (rule of thumb: use  $r_{ap}^{(1)} = 1.5 \times \text{FWHM}$ ).
- Select a *bright* star in the field and measure the counts  $N_{ap1}$  inside  $r_{ap}^{(1)}$  and the counts  $N_{ap2}$  inside a “big” aperture  $r_{ap}^{(2)}$ .
- Aperture correction  $c(\text{ap}) = N_{ap2}/N_{ap1}$ .
- Corrected counts  $N_{apcorr} = c(\text{ap}) \times N$
- Works for point sources only!

## 1.2 PSF photometry (point sources only)

PSF = “Point Spread Function” i.e. how a point source looks through the optical system.

Stars are practically point sources  $\rightarrow$  PSF = image of a star.

PSF photometry :

1) Create the PSF model = 2-dim image of a point source

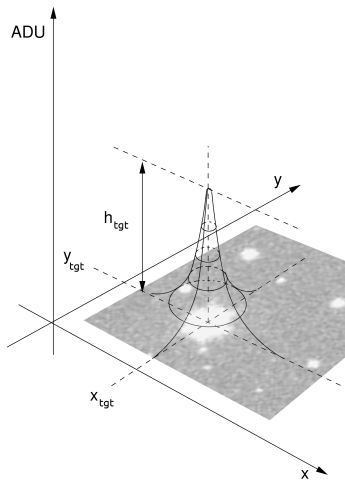
- Can use a theoretical model, an analytic function or an image of a sufficiently bright star in the CCD frame (the last option usually gives the best results).
- Common analytical functions:

$$\text{Gauss : } I(r) = I_0 \exp(-r^2/2\sigma^2)$$

$$\text{Moffat : } I(r) = I_0 \left[ 1 + K \left( \frac{r}{R} \right)^2 \right]^{-\beta}, \quad R = \text{FWHM}/2, \quad K = 2^{1/\beta} - 1$$

# PSF photometry

- 2) Obtain aperture photometry of your PSF  $\rightarrow$  relationship between total counts and height :  $N_{PSF} = k * h_{PSF}$
- 3) Fit the PSF model to the target. Free parameters of the fit :  $x_{tgt}$ ,  $y_{tgt}$ ,  $h_{tgt}$ .
- 4)  $N_{tgt} = k * h_{tgt}$



## 1.3 Sky estimation

- Select the sky annulus as close to the target you can without introducing target light into the annulus.
- Check that there are no bright targets in the annulus.
  - IRAF can cope to some extent with stars in the sky annulus by using sigma clipping and by using mode of the sky pixel distribution as  $N_{sky}$ . In practice IRAF computes the mode from

$$mode = 3 \times median - 2 \times mean$$

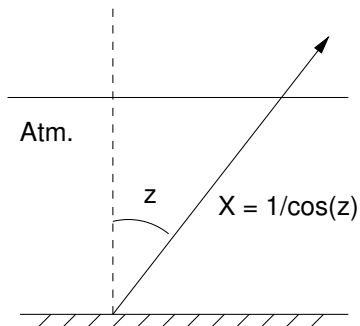
- Sky estimation is important for faint targets, where it can easily be the largest source of error.

## 2. Correcting the counts



## 2.1 Atmosphere

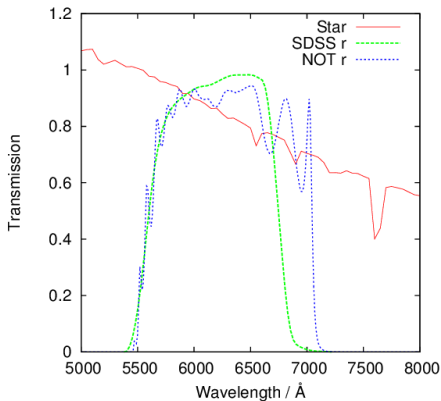
- Absorbs light in an altitude-dependent way.



- Parameterized through *airmass*  $X$ .  
 $z = 0 \text{ deg} \rightarrow X = 1$   
 $z = 60 \text{ deg} \rightarrow X = 2$

## 2.2 Instrumental effects

- Bias, dark, “cosmic rays” create artificial signal.
- Pixels have different gains.  
→ Corrected by image processing.



- Filter mismatch: e.g. “r-band flux” is a *monochromatic* quantity [ $\text{erg s}^{-1} \text{cm}^{-2} \text{Å}^{-1}$ ] defined at some wavelength.

However, it is measured with a broad filter → the result depends on the transmission curve of the filter.

### 3. Calibration

## 3.1 Magnitudes

Traditionally, the calibration is done using magnitudes. The relationship between fluxes  $F$  and magnitudes  $m$  is

$$F = F_0 10^{-m/2.5}, \quad (3)$$

where  $F_0$  is the flux corresponding to magnitude 0. By taking the logarithm and multiplying by -2.5 we get

$$m = 2.5 \log_{10} F_0 - 2.5 \log_{10} F = m_0 - 2.5 \log_{10} F. \quad (4)$$

If we observe  $N_s$  counts/s from a target, we can write

$$F = c N_s, \quad (5)$$

where  $c$  is a constant. Thus

$$m = \underbrace{m_0 - 2.5 \log_{10} c}_{zp} - 2.5 \log_{10} N_s. \quad (6)$$

## 3.2 Calibration equation

$$m_c = zp - 2.5 \log N - k \times X - C \times Color + 2.5 \log T_{exp}, \quad (7)$$

where

$N$  = observed counts

$X$  = airmass

$Color$  = difference of target magnitude in two filters, e.g. (B-V)

$T_{exp}$  = exposure time

$zp$  = system zero point (drifts slowly with time)

$k$  = extinction coefficient (varies from night to night)

$C$  = color term of the filter (does no change).

Calibration = determine  $zp$ ,  $k$  and  $C$ .

# Absolute calibration

- You have obtained CCD-images of several calibration targets (“standard stars”) with known  $m$  and  $color$ . The goal is to determine  $z_p$ ,  $k$  and  $C$  from these observations.
- This method works only if the night is *photometric*, i.e. extinction  $k$  stays constant.
- It is important to cover a wide range of airmasses (from 1 to 2) and colors.

## Absolute calibration : $zp$ , $k$ and $C$ .

- Measure  $N$  for each standard star.
- Compute *instrumental magnitudes*  $m_i$  from

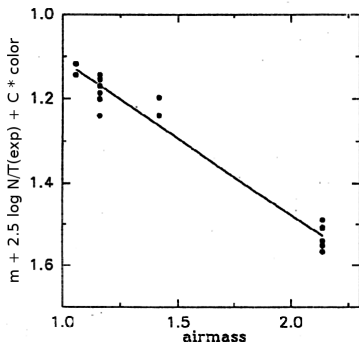
$$m_i = m_0 - 2.5 \log_{10} N + 2.5 \log_{10} T_{exp} , \quad (8)$$

where  $m_0$  is an arbitrary zero point (IRAF default 25.0).

- Plot each observation in a 3-d space with  $X$ , *color* and  $m - m_i$  as axes, where  $m$  is the catalog magnitude of the star. The points should fall on a plane (Eq. 7).
- Fitting a plane to this 3-d data will give you  $zp$ ,  $k$  and  $C$ .

# Absolute calibration : $z_p$ , $k$

- Usually  $C$  is known for the system and it doesn't vary. It is then sufficient to determine only  $z_p$  and  $k$  from a 2-d fit.
- For each standard star observation, plot  $X$  in the x-axis and  $m - m_i + C \times color$  in the y-axis. A linear fit will then yield  $z_p$  and  $k$ .





- Assume that we have on the CCD image the target and a another target with known magnitude  $m_{ref}$ .
- If we mark with  $m$  the magnitude of our target of interest, we get from Eq. (7)

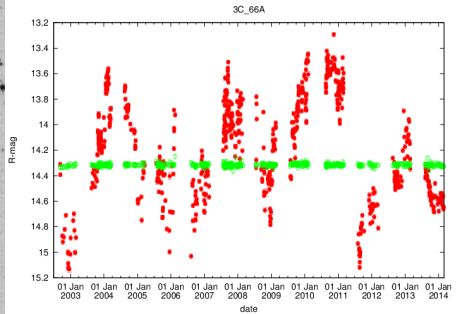
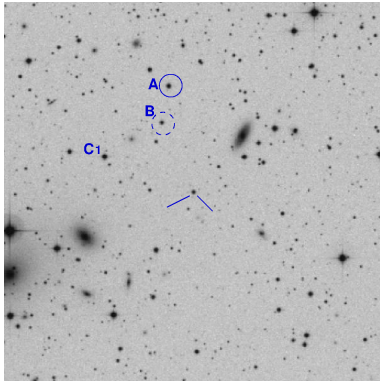
$$m = m_{ref} + 2.5 \log \frac{N_{ref}}{N} - C \times (Color - Color_{ref}) \quad (9)$$

→  $m$  can be easily computed if  $C$  is known.

# Advantages of differential calibration

- The zero point and atmospheric effects are almost completely eliminated. In practice small residual effects persist at  $\sim 0.01$ - $0.02$  mag level and below.
- Works well even through clouds.
- No need for calibration observations if good quality survey data exist for the observed field and filters (e.g. SDSS).
- No need for aperture correction.

# Tuorla blazar monitoring



# From magnitudes to fluxes

$$F = F_0 * 10^{-0.4*m_c} , \quad (10)$$

where  $F_0$  is a filter-dependent zero point, for instance:

Filter	$F_0$ [Jy]
U	1790
B	4263
V	3636
R	3064
I	2416
SDSS u,g,r,i,z	3631